

UNIVERSITY OF SURREY[©]

**School of Electronics and Physical Sciences
Department of Mathematics**

M.Math Undergraduate Programme in Mathematical Studies

Examination

Module MSM.TWW THEORY OF WATER WAVES

Time allowed - 2.5 hrs

Spring Semester 2007

Attempt FOUR questions.

If any candidate attempts more than FOUR questions,
only the best FOUR solutions will be taken into account.

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Question 1

(Quintic nonlinear Schrödinger equation)

The nonlinear Schrödinger equation with quintic nonlinearity, for the complex-valued function $A(x, t)$, is

$$i A_t + A_{xx} + a |A|^4 A = 0, \quad a = \pm 1. \quad (1)$$

Associated with quintic NLS are the conserved densities

$$M = |A|^2, \quad I = -i(\bar{A}A_x - A\bar{A}_x), \quad E = |A_x|^2 - \frac{a}{3}|A|^6.$$

(a) Find fluxes Q , S and F such that when A is a solution of (1) then

$$M_t + Q_x = 0, \quad I_t + S_x = 0, \quad E_t + F_x = 0.$$

[6]

(b) Suppose $A(x, t) \rightarrow 0$ as $x \rightarrow \pm\infty$. Show that

$$\int_{-\infty}^{+\infty} S(x, t) dx = C \int_{-\infty}^{+\infty} E(x, t) dx,$$

and determine the value of the constant C .

[5]

(c) Suppose $A(x, t) \rightarrow 0$ as $x \rightarrow \pm\infty$ and consider the moment of inertia of a solution,

$$J(t) = \int_{-\infty}^{+\infty} x^2 |A(x, t)|^2 dx.$$

By differentiating $J(t)$ and using the conservation laws show that

$$\frac{d^2 J}{dt^2} = 2 \int_{-\infty}^{+\infty} S(x, t) dx.$$

[9]

(d) What can one conclude about existence of solutions of (1)? Consider the cases $a = -1$ and $a = +1$ separately.

[5]

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Question 2

(Conservation of wave action)

Consider the nonlinear wave equation

$$u_{tt} - u_{xx} + u = \varepsilon u^3, \quad (2)$$

where ε is a small parameter.

- (a) For the linear equation, $u_{tt} - u_{xx} + u = 0$, determine the dispersion relation and group velocity for normal-mode solutions: $u(x, t) = Ae^{i(kx - \omega t)} + c.c.$ [3]
- (b) Let $X = \varepsilon x$ and $T = \varepsilon t$ and define a phase by $\theta_x = k(X, T)$ and $\theta_t = -\omega(X, T)$. Using the chain rule:

$$\frac{\partial}{\partial t} = -\omega \frac{\partial}{\partial \theta} + \varepsilon \frac{\partial}{\partial T}, \quad \frac{\partial}{\partial x} = k \frac{\partial}{\partial \theta} + \varepsilon \frac{\partial}{\partial X},$$

transform the PDE into an equation for $u(\theta, X, T, \varepsilon)$. [4]

- (c) Let

$$u(\theta, X, T, \varepsilon) = u_0(\theta, X, T) + \varepsilon u_1(\theta, X, T) + \dots,$$

and take $u_0 = A(X, T)e^{i\theta} + \overline{A(X, T)}e^{-i\theta}$. By requiring u_1 to be 2π -periodic in θ , show that the solvability condition requires $A(X, T)$ to satisfy

$$\omega_T A + 2\omega A_T + k_X A + 2k A_X - 3i|A|^2 A = 0. \quad (3)$$

[10]

- (d) Show that *conservation of wave action*

$$\frac{\partial}{\partial T} \left(\frac{E}{\omega} \right) + \frac{\partial}{\partial X} \left(c_g \frac{E}{\omega} \right) = 0, \quad \text{where } E = \omega^2 |A|^2,$$

follows from (3).

[8]

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Question 3

(Wave refraction and ray theory)

Consider the steady propagation of waves in shallow water governed by the ray equations. The graph of a ray is given by $(X, Y(X))$ with $Y(X)$ satisfying

$$\frac{d}{dX} \left(\frac{\sigma Y_X}{\sqrt{(1+Y_X^2)}} \right) - \sigma_Y \sqrt{(1+Y_X^2)} = 0. \quad (4)$$

where $\sigma(X, Y)$ is determined from the equation

$$g\sigma(X, Y) \tanh(\sigma(X, Y)D(X, Y)) = \text{constant},$$

and $D(X, Y)$ is determined by bottom topography. You may assume throughout that $\frac{dY}{dX} > 0$.

(a) Suppose $\sigma_Y = 0$ and $\sigma'(X) > 0$ (signifying decreasing depth). Show that

$$\frac{d}{dX} \left(\frac{dY}{dX} \right) < 0.$$

[5]

(b) Let $Y_X = \tan(\theta)$ with $0 < \theta < \pi/2$ and $\sigma_Y = 0$. Show that rays satisfy Snell's Law

$$\sigma(X) \sin\theta = \text{constant}.$$

[6]

(c) Suppose $\sigma(Y)$; that is, a function of Y only. Show that rays satisfy

$$\frac{\sigma(Y)}{\sqrt{1+Y_X^2}} = \text{constant}.$$

[8]

(d) Suppose as in part (c) that σ depends on Y only, and $\sigma_Y > 0$. Does

$$\frac{d}{dX} \left(\frac{dY}{dX} \right)$$

increase or decrease along rays?

[6]

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Question 4

(The Korteweg-DeVries equation)

Consider the KdV equation with variable dispersion

$$u_t + uu_x + a(x)u_{xxx} = 0. \quad (5)$$

- (a) Suppose $a(x) = -1$ for all x (negative dispersion) and consider travelling wave solutions of (5) of the form $u(x, t) = \eta(\xi)$ with $\xi = x + ct$ with positive wave speed $c > 0$. Show that η satisfies

$$-\eta_{\xi\xi} + \frac{1}{2}\eta^2 + c\eta = \text{constant}. \quad (6)$$

[3]

- (b) Take $a = -1$, $c > 0$ and $\text{constant} = 0$. Show that there is a solitary wave solution $\eta(\xi)$ of (6) with $\eta(\xi) \rightarrow 0$ as $\xi \rightarrow \pm\infty$. Give an explicit expression for $\eta(\xi)$. How is this solitary wave different from the classical KdV solitary wave where $a(x) := +1$?

[8]

- (c) Consider (5) with $a(x)$ a given smooth function which is not identically zero and let $M(x, t) := u(x, t)$ and $Q(x, t) = \frac{1}{2}u^2 + a(x)u_{xx}$. Show that

$$M_t + Q_x = R_1(x, t),$$

where $R_1(x, t)$ is *not zero* when $a(x)$ is nonconstant, even when $u(x, t)$ satisfies (5). Give an expression for $R_1(x, t)$.

[6]

- (d) Similar to part (c) show that the momentum conservation law for (5) is perturbed to

$$I_t + S_x = R_2(x, t), \quad I(x, t) = \frac{1}{2}u(x, t)^2,$$

when $u(x, t)$ is a solution of (5). Give expressions for $S(x, t)$ and $R_2(x, t)$.

[8]

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Question 5

(Dispersion properties of interfacial waves)

Consider a fluid configuration with two layers of fluid with different densities and irrotational flow. The lower layer has density ρ_1 and extends from $y = 0$ to $y = h$. The upper layer has density ρ_2 and extends from $y = h$ to $y = +\infty$.

The governing equation and boundary conditions are

$$\begin{aligned} \Delta\phi_1 = 0 & \quad \text{for } 0 < y < h, & \quad \frac{\partial\phi_1}{\partial y} = 0 \text{ at } y = 0, \\ \Delta\phi_2 = 0 & \quad \text{for } h < y < +\infty, & \quad \frac{\partial\phi_2}{\partial y} \rightarrow 0 \text{ as } y \rightarrow +\infty, \end{aligned}$$

where Δ is the Laplacian. The boundary conditions at the interface $y = h$ are

$$\frac{\partial\eta}{\partial t} = \frac{\partial\phi_1}{\partial y}, \quad \frac{\partial\phi_1}{\partial y} = \frac{\partial\phi_2}{\partial y}, \quad \rho_1 \frac{\partial\phi_1}{\partial t} - \rho_2 \frac{\partial\phi_2}{\partial t} + (\rho_1 - \rho_2)g\eta = 0,$$

where g is the positive gravitational constant.

- (a) Consider normal mode solutions of the form $\eta(x, t) = Ae^{i(kx - \omega t)} + c.c.$ and

$$\phi_1(x, y, t) = B_1(y)e^{i(kx - \omega t)} + c.c. \quad \text{and} \quad \phi_2(x, y, t) = B_2(y)e^{i(kx - \omega t)} + c.c.,$$

where A is a complex constant. Determine expressions for $B_1(y)$ and $B_2(y)$ satisfying the Laplace equation in interior and the boundary conditions at $y = 0$ and $y \rightarrow \infty$. [6]

- (b) Using the boundary conditions at $y = h$ determine a relationship between $B_1(h)$ and A and $B_2(h)$ and A . [7]

- (c) Show that the dispersion relation for the system is

$$\omega^2 = \frac{(\rho_1 - \rho_2)gk \tanh(kh)}{\rho_1 + \rho_2 \tanh(kh)}. \quad [8]$$

- (d) Suppose $k > 0$ and $\rho_1 < \rho_2$. What does this mean physically? Using the dispersion relation as a guide, discuss the implication for the time evolution. [4]

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